

# Investigating protein dynamics in collective coordinate space

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Currently, collective coordinates are commonly employed in order to examine protein dynamics. In recent studies, they have been successfully applied to finding functionally relevant motions, to investigating the physical nature of protein dynamics, to sampling of the conformational space and to the analysis of experimental data. Collective coordinates also have other possible applications.

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## Abbreviations

<b>BPTI</b>	bovine pancreatic trypsin inhibitor
<b>ED</b>	essential dynamics
<b>JAM</b>	jumping-among-minima
<b>MC</b>	Monte Carlo
<b>MD</b>	molecular dynamics
<b>MODC</b>	molecule optimal dynamic coordinate
<b>MSF</b>	mean square fluctuation
<b>NM</b>	normal mode
<b>NMA</b>	NM analysis
<b>PC</b>	principal component
<b>PCA</b>	PC analysis

## Introduction

Dynamics is essential for a protein to function. Recent advancements in computers and computational techniques enable us to simulate protein dynamics more accurately in a more realistic environment; however, it is not a trivial problem to extract a functionally relevant motion from simulation results. A key solution to this problem is the use of collective coordinates. Recent studies have shown that functionally relevant motions occur along the direction of a few collective coordinates, which dominantly contribute to atomic fluctuation. One example is the molecular dynamics (MD) simulation of CRP:(cAMP)<sub>2</sub>, the cAMP receptor protein bound to cAMP, carried out by García and Harman [1]. Two conformations, open and closed, exist in the crystal environment, depending on the occurrence of binding to DNA. The transition from the open to the closed conformation found during a MD simulation involves only two collective coordinates. Another example is the case of the T4 lysozyme [2•], whereby a large variation of the hinge-bending angle has been seen in the crystal environment. The first two largest amplitude collective modes determined from the MD trajectory, which describes the domain motion, are identified as being the closure and twisting modes.

More importantly, collective coordinates can be used to predict a low-dimensional subspace in which significant

protein motion is expected to take place. If this concept is true, the number of degrees of freedom can be extremely reduced. In a Cartesian coordinate space protein, the number of internal degrees of freedom is  $(3N-6)$ , where  $N$  is the number of atoms considered. In the case of lysozyme, where  $N \cong 2000$ , the number of degrees of freedom can be reduced to 5% if 300 appropriately chosen collective coordinates span the subspace. It is well known that atomic fluctuations within such a subspace dominantly contribute to total fluctuation. A subspace that is spanned by a small number of low-frequency normal modes is termed an ‘important subspace’ [3]. This concept has been successfully employed to refine crystal structures [3–5]. Later, Amadei *et al.* [6] termed a subspace that is spanned by a subset of anharmonic collective coordinates determined by MD an ‘essential subspace’. A small number of large amplitude collective modes determined by simulations depends on the simulation length in some cases [7], because large amplitude transitions occur stochastically. Teeter and Case [8] showed that the very low frequency modes are sensitive to the force-field; however, the frequency distributions and fluctuations were remarkably similar. In addition, the subspace spanned by a small number of collective coordinates is invariant in many cases and is almost independent of the treatment of the degrees of freedom [9,10], the solvent effect [11,12], the simulation length [13••] and the initial and other conditions [14]. In this paper, we review how the concept of these subspaces is a powerful tool for investigating protein dynamics.

## Determining a collective coordinate set

A collective coordinate set can be determined as a set of eigenvectors that diagonalizes a second moment matrix. This is based on what is called principal component analysis (PCA). In Cartesian coordinate space, an element of the second moment matrix  $\mathbf{A}$ ,  $a_{ij}$ , is defined as:

$$a_{ij} \equiv \langle (x_i - \langle x_i \rangle)(x_j - \langle x_j \rangle) \rangle \quad (1)$$

where  $\langle \dots \rangle$  is the average over the instantaneous structures sampled during the period of simulation and  $x_i$  is a mass-weighted atomic coordinate. A set of eigenvalues and eigenvectors is obtained by solving the standard eigenvalue problem:

$$\mathbf{A}\mathbf{W} = \mathbf{W}\boldsymbol{\zeta} \quad (2)$$

where  $\boldsymbol{\zeta} = \text{diag}(\zeta_m)$  is a diagonal matrix whose  $m$ th diagonal element is the eigenvalue  $\zeta_m$ ,  $\mathbf{W}$  is a matrix whose  $m$ th column vector  $w_m$  is the eigenvector of  $\zeta_m$ . The vector  $w_m$ , given by a linear combination of Cartesian coordinates, represents the axis of the  $m$ th collective coordinate in the conformational space. The eigenvalue  $\zeta_m$  represents the

mean square fluctuation (MSF) along this axis. The second moment matrix is calculated directly from a MD or Monte Carlo (MC) trajectory. The motions along these coordinates are termed 'modes', although they are not necessarily vibrational. In normal mode analysis (NMA), the second moment matrix is given by:

$$\mathbf{A} = k_B T \mathbf{F}^{-1} \quad (3)$$

where  $\mathbf{F}$  is a matrix of second derivatives of the potential energy at a minimum point,  $k_B$  is the Boltzmann constant and  $T$  is the absolute temperature. In dihedral angle space, the number of degrees of freedom is reduced to about an eighth. For established methods in dihedral angle space, see the review by Hayward and Go [15].

Further explanation may be required concerning the term 'PCA'. When PCA was first applied to a MD trajectory, it was termed 'quasi-harmonic', because the focus was on the anharmonic aspects of protein dynamics [16]. The term 'principal component analysis' was first used in [11] because the projection of the MD trajectory onto a subspace was more highlighted. The term 'essential dynamics' (ED) is also used to represent a few anharmonic collective modes [6]. García and Harman [1] called principal coordinates 'molecule optimal dynamic coordinates' (MODCs). These analyses share essentially the same concept, apart from some details, for example, the atoms considered, whether it is mass weighted or not and so on. We prefer the use of mass-weighted coordinates because the relationship between PCA and NMA is clearer.

Practically, the most time consuming aspect of PCA and NMA is the diagonalization of the second moment matrix. In PCA, the reduction of the matrix can be done easily by selecting solely heavy atoms or C $\alpha$  atoms. As such a reduction cannot be done, diagonalization is more involved in NMA. For recent advancements in both NMA and diagonalization techniques, see the review by Hayward [17]. Simplified models have been proposed and have been successful in predicting atomic fluctuations [18,19\*\*]. These methods are designed to reproduce a few large-amplitude collective modes using simplified force-fields.

As an alternative to the second moment matrix, the metric matrix has been employed to determine collective coordinates in a method termed 'multidimensional scaling' [20].

### Projection of protein motion onto collective coordinate space

One of the first attempts to project MD and MC trajectories onto the normal mode (NM) space was carried out on the human lysozyme [21]. This was successful in explaining the anharmonicity in a single conformational substate. Projection onto the NM space is not, however, an efficient way of observing intersubstate transitions [22]. Two important features, transitions between substates and solvent effects, were clearly revealed by the first attempt at

projecting a MD trajectory onto the principal coordinate space [9]. By projecting a MD trajectory of melittin in water onto the two-dimensional space spanned by the largest and second largest amplitude principal modes, a transition from one substate to another is very clearly seen. The motion along these coordinates involves a bending of the molecule. Soon after, conformational transitions in the subspace spanned by a few principal components (PCs) were reported in the cases of Met-enkephaline [23], crambin [24], bovine pancreatic trypsin inhibitor (BPTI) [12] and hen-egg lysozyme [6]. PCA has now become a common tool for investigating important motions of biological molecules [25–29].




In PCA, solvent effects are evident in two ways. One is detected as small-amplitude conformational transitions that are not seen in a vacuum [11,12]. In the case of melittin, this type of transition involves a change of hydration structure [30]. Fine structures in the free energy surface, which are not seen in a vacuum, exist in water [31]. In the case of myoglobin, global transitions were observed more frequently in a fully solvated model, compared with those in a partially solvated model [32]. This may be due to the fact that the presence of solvent lowers the barriers of large conformational transitions. The other solvent effect appears as a damping of collective motion. In the case of melittin [11] and BPTI [12], time correlation functions of the large-amplitude PCs show overdamping behavior. In the case of a harmonic energy surface, this damping effect can be elegantly incorporated into NMA. 'Langevin mode' analysis was first proposed by Lamm and Szabo [33] and was applied to crambin and DNA by Kottalam and Case [34]. The overdamping and underdamping behaviors of collective modes were elucidated by this method; however, the problem as to how to estimate the matrix of friction coefficients remains. From a back calculation of this matrix from an MD trajectory, it has been shown that simple hydrodynamic models for the friction matrix fail to accurately predict off-diagonal elements of the friction matrix [11]. If friction coefficients are correctly given, Langevin mode analysis can explain the spectral shape determined by an inelastic neutron scattering experiment in a low-frequency region [12].

It should be also noted that there is an important feature that is not affected by solvent. Using the cross-correlations and co-variances of atomic fluctuations, Ichiye and Karplus [35] showed that the correlations of internal motions are similar in different environments. This is consistent with the finding that a subspace spanned by a small number of collective coordinates is almost invariant, as described above.

By projecting an MD trajectory onto principal coordinates, Caves *et al.* [36\*\*] reported a problem of sampling using MD. Using 10 independent MD trajectories of crambin with different initial velocities, they found four major conformational regions. In this case, the use of multiple trajectories is more efficient in sampling compared to a single, long trajectory. An alternative way of solving the

Table 1

## Characteristics of collective modes determined by PCA and JAM\*.

	Multiply hierarchical mode	Singly hierarchical mode	Harmonic mode
Free energy surface			
Probability distribution	Non-Gaussian multiple peak	Gaussian-like single peak	Gaussian
Anharmonicity factor $\rho_m^\dagger$	> 2	2-1	= 1
Packing topology	Significant change	No change	No change
Fluctuation of dihedral angles	Large change in both mainchains and sidechains	Large change in sidechains	Small change
Possible relationship with other classifications	Essential subspace [6] Relaxation dynamics [1]	Fluctuation dynamics [1]	Unimodal dynamics [1]
	Equivalent to NM important subspace [3]		
<b>Specific data for the human lysozyme</b>			
Principal mode number	1-30	31-300	301-6117
% in mode number	0.5	4.5	95.0
% in total MSF	82.1	15.0	2.9
Transition timescale (s)	> 10 <sup>-13</sup> (various transition timescales co-exist in one mode)	10 <sup>-13</sup> -10 <sup>-11</sup>	-
RMSF (Å)	> 0.08	0.08-0.015	0.015 >

\*A newly generated table based on [13\*\*]. <sup>†</sup>The factor  $\rho_m$  is defined as  $\rho_m = \left( \zeta_m / \zeta_m^{har} \right)^{1/2}$ , where  $\zeta_m^{har}$  is the expected MSF from NM along the direction [6]. RMSF, root mean square fluctuation.

sampling problem is the use of efficient sampling methods, for example, multicanonical MC [37] and MD [38]; however, these still need enormous amounts of computational time, even in vacuum [39]. Recently, Becker [40] examined the possibility of using PCA for the mapping of the potential energy surface of small peptides.

### Energy landscapes in collective coordinate space

After pioneering work on the energy landscapes of myoglobin [41] and BPTI [42-46] in their native states, most of the recent attempts to investigate energy landscapes have been carried out using collective coordinates. From non-Gaussian probability distribution functions of projections of an MD trajectory onto principal coordinates, large-amplitude nonlinear motions were identified [24]. Similarly, Amadei *et al.* [6] termed the space spanned by these anharmonic coordinates 'essential subspace'. In order to further understand harmonic and anharmonic aspects of protein dynamics, NMs have been employed [22,47]. Using what is called the 'anharmonicity factor', anharmonic modes and harmonic modes are clearly separated by principal mode numbers [22]. In the case of BPTI, only 11.7% of the modes are anharmonic (equivalent to 7.4% in an all-atom model) and, in the case of human lysozyme [13\*\*], only 5.0% are anharmonic. By comparing the NMs of 201 different minima, Janežič *et al.* [48] found that low-frequency NMs extensively overlap in angular frequency when minima are separated by a relatively short time interval. Using

different time windows, Troyer and Cohen [20] found different types of conformational transitions. Transitions among seven major clusters in the binding loop of BPTI were suggested to be functionally important as the inhibitor. García and Harman [1] pointed out three types of dynamics — unimodal, fluctuation and relaxation. In the case of crambin in a crystal environment, a motion whose characteristic time is slower than 2 ns was found [49\*\*].

The jumping-among-minima (JAM) model elucidates the above findings more clearly [13\*\*]. In this model the second moment matrix is divided into two terms:

$$a_{ij} = \sum_{k=1}^M f_k \left( \langle x_i \rangle_k - \langle x_i \rangle \right) \left( \langle x_j \rangle_k - \langle x_j \rangle \right) + \sum_{k=1}^M f_k \left( \langle x_i - \langle x_i \rangle \rangle_k \right) \left( \langle x_j - \langle x_j \rangle \rangle_k \right) \quad (4)$$

where  $M$  is the number of conformational substates,  $k$  is the index of the substate,  $\langle \dots \rangle_k$  is an average over the period during which the state point stays in the  $k$ th substate and  $f_k$  is a fraction of this period to the total simulation time. The first and second terms of Equation 4 represent contributions from atomic fluctuations arising from JAM and arising from fluctuations within each substate, respectively. This model was applied to the analysis of a 1 ns MD trajectory of the human lysozyme in water. From the analysis of intrasubstate motion,

each conformational subspace is shown to be nearly harmonic and mutually similar. Using the anharmonicity factor and analyzing the shapes of the probability distribution functions, three types of modes are identified: multiply hierarchical modes, singly hierarchical modes and harmonic modes. Along the axis of multiply hierarchical modes, different levels of energy barriers co-exist; see Table 1 for details.

Beyond the native state, we think that collective coordinate descriptions of protein dynamics would still be useful. A test calculation shows that, even in the wider conformational space beyond the native state, a small number of PCs still dominantly contribute to the total fluctuation (A Kitao, H Hirose, S Nakamura, M Ikeguchi, K Shimizu, N Go, unpublished data). This suggests the possibility of using PCA to examine the hypotheses of the folding pathway and the folding funnel by projecting the trajectories of a folding/unfolding simulation onto a low-dimensional subspace.

### Functional importance of collective motion

As already described above, PCA is a powerful tool for determining functionally relevant motions. Recently, in order to further understand the characteristics of collective motions, the concept of a 'dynamic domain' has been introduced, whereby collective motion is reduced to rigid-body domain motion [2\*,50,51,52\*\*]. Using well-defined criteria, dynamic domains and rotation axes have been determined. The domain motion is quantified by a rotation angle around the axis. The concept of a dynamic domain is expected to give some insights into understanding protein structure.

There are some cases for which the functional importance of collective motion is not evident from the motion itself. Using NM variables, an analytical model for formulating a protein electron transfer reorganization energy spectrum was proposed [53\*] and applied to cytochrome c [54\*]. This is an extension of the work done by Schulten and Tsch [55] to include a collective coordinate description so that the structural aspects of the phenomenon can be discussed explicitly. This model enables the identification of NMs that contribute significantly to the reorganization energy. As described in the original studies [53\*,54\*], it can be easily extended to a more realistic model using PCA.

### Using collective coordinate as a basis set for efficient sampling

As protein dynamics is highly anisotropic, MC simulations with isotropic sampling step sizes are inefficient. In order to overcome this, NMs are used as variables and the sampling step size is scaled by NM amplitudes [56]. In MD, collective coordinates that span essential space are used as variables [57]. These methods remarkably improve the sampling efficiency, although the problem still remains as to how the solvent effect is incorporated into the former method. One possible way is the use of the XRISM theory. Although extensive calculation is needed, XRISM has been applied to a 436 site (atom) model of melittin [31]. A recent,

efficient algorithm [58] could be incorporated into these efficient sampling methods.

### Analysis of experimental data by means of collective variables

In B-factor refinement of X-ray crystal structures, atomic motion is assumed to be isotropic. Although anisotropic B-factor refinement has been tried, it requires that more parameters be determined. Using the NM important subspace, the number of parameters is highly reduced in NM refinement [3–5,59]. In addition, external and internal motions can be clearly separated. The internal motion of the human lysozyme in its crystal environment is highly anisotropic and atomic fluctuations are in good agreement with those determined by NM refinement [3]. The refinement of the NMR spectroscopy structure was also carried out using the same concept [60]. Using PCA, Abseher *et al.* [61] found a significant overlap between those conformations sampled by unconstrained MD and those determined in order to satisfy NMR constraints. This is as expected from the concept discussed in the Introduction. Collective coordinates have also been used to estimate NMR order parameters [62,63].

The JAM model can also be used to analyze experimental data obtained by X-ray and neutron scattering, laser spectroscopy and NMR spectroscopy. If intrastate motion is given by NM and is invariant in all conformation substates, the atomic coordinate  $x_i(t)$  is given by the sum of the contributions from external motion  $x_i^{EXT}(t)$ , JAM motion  $x_i^{JAM}(t)$  and NM vibration  $x_i^{NM}(t)$  as:

$$x_i(t) = x_i^{EXT}(t) + x_i^{JAM}(t) + x_i^{NM}(t). \quad (5)$$

For example, in the case of inelastic neutron scattering, the incoherent dynamic structure factor  $S_{inc}(Q, \omega)$  is given by the convolution of three structure factors —  $S_{inc}^{EXT}(Q, \omega)$ ,  $S_{inc}^{JAM}(Q, \omega)$  and  $S_{inc}^{NM}(Q, \omega)$  as:

$$S_{inc}(Q, \omega) = S_{inc}^{EXT}(Q, \omega) \otimes S_{inc}^{JAM}(Q, \omega) \otimes S_{inc}^{NM}(Q, \omega) \quad (6)$$

if these three types of motion are uncorrelated. Here, the symbol  $\otimes$  denotes the convolution product over angular frequency. In the case of staphylococcal nuclease, the structure factor measured by TFXA was in remarkably good agreement with the calculated  $S_{inc}^{NM}(Q, \omega)$  factor, except in a few high-frequency peaks and very low-frequency regions [64\*\*]. It may also be possible to characterize the JAM motion from a high-resolution structure factor.

### Conclusions

Collective coordinates are employed to investigate protein dynamics. They are useful not only in finding functionally relevant motions, but also in understanding the physical nature of the protein energy landscape, efficient sampling and the analysis of experimental data.

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